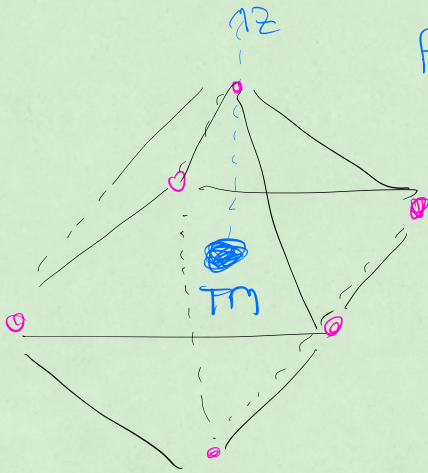


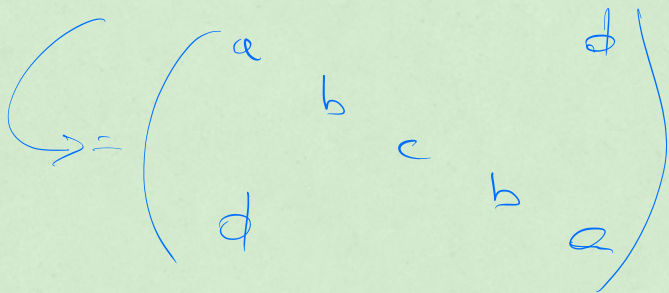
17.10.2022



Perfect octahed.

$O_h$

$$\langle l' m' | H_{CF} | l m \rangle$$



Diagonalization

eigenstates

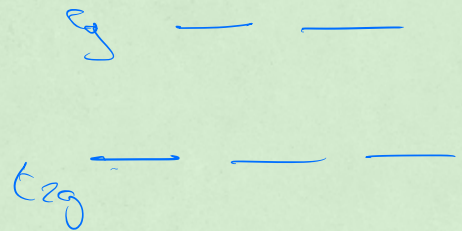
$|d_{...}\rangle$

Symmetry adapted w.f.

$O_h$



into SOLID



$$H = H_0 + H_{ee} + H_{CF} + \dots$$

+ Pauli p.p.  
e.g. 3d

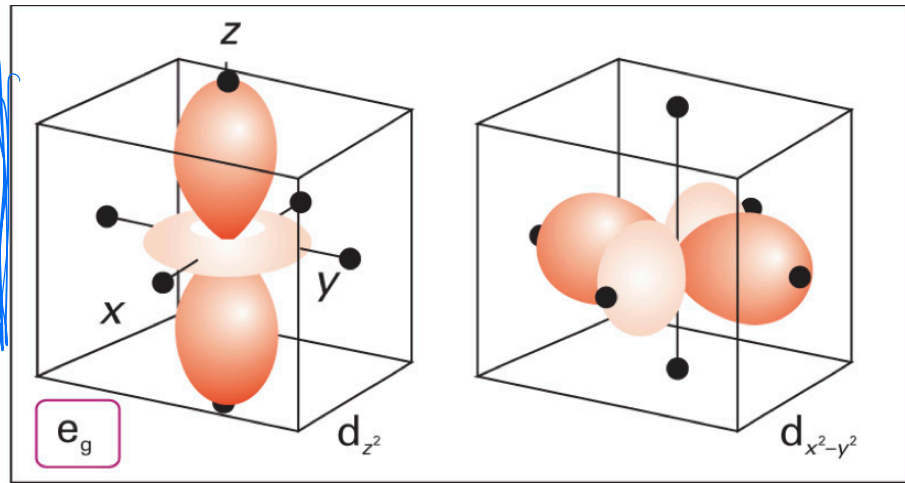
$t_{2g}, e_g$

$|d_{z^2}\rangle, \dots$   
Single el. levels

single el. levels

$$|l, m\rangle \approx Y_{l, m}(\Omega)$$

elongation



$$d_{x^2-y^2} = \frac{1}{\sqrt{2}} [ |2, +2\rangle + |2, -2\rangle ]$$

$$d_{z^2} = |2, 0\rangle$$

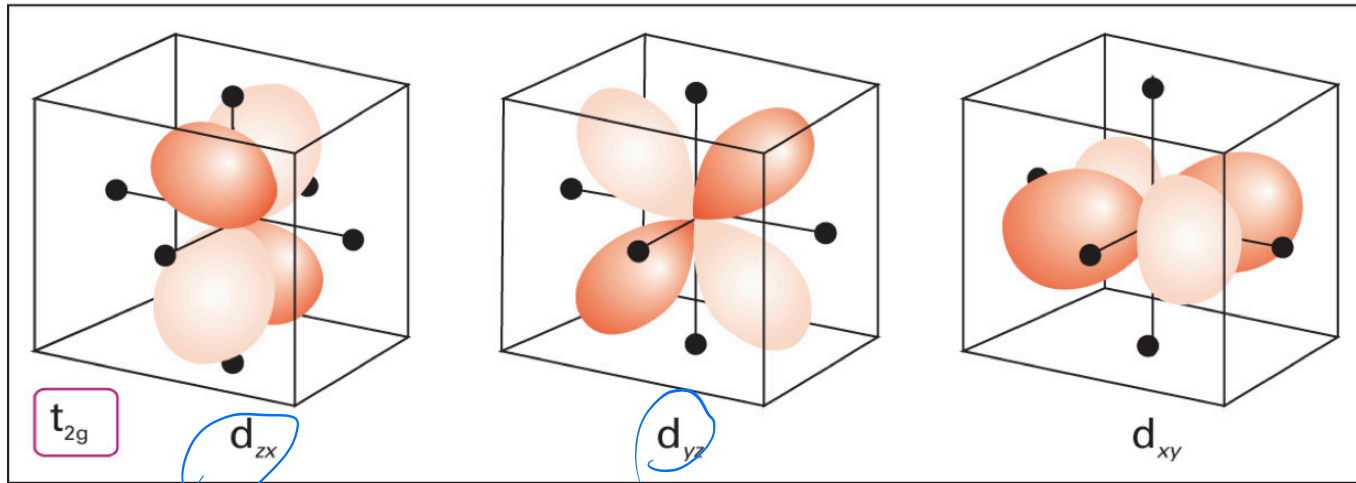
}  $e_g$

$$d_{xy} = -\frac{i}{\sqrt{2}} [ |2, +2\rangle - |2, -2\rangle ]$$

$$d_{xz} = -\frac{1}{\sqrt{2}} [ |2, +1\rangle - |2, -1\rangle ]$$

$$d_{yz} = \frac{i}{\sqrt{2}} [ |2, +1\rangle + |2, -1\rangle ]$$

}  $t_{2g}$



# NOMENCLATURE

	Single el. level	Multi el. level
free ion	$(3d)^{n_e}$	${}^2P_{1/2} \dots$ (Russel Saunders terms)
CF- $O_h$	$t_{2g}, e_g$	${}^{2S+1}A_{\dots}, {}^{2S+1}E, T$

Transition metals

	A	E	T
orbital deg	1	2	3

J<sub>intra</sub>

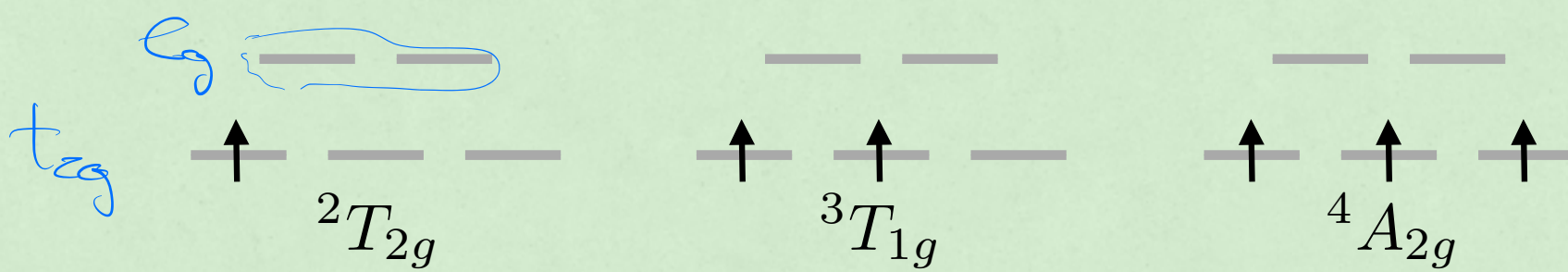
He excited states  $(1s)^1 (2s)^1$

SPATIAL $\psi^T(r_1, r_2)$	Asym. $1 \leftrightarrow 2$	Spin w.f. $\chi^T$
$\psi^S(r_1, r_2)$	Symm. $2 \leftrightarrow 1$	$\chi^S$

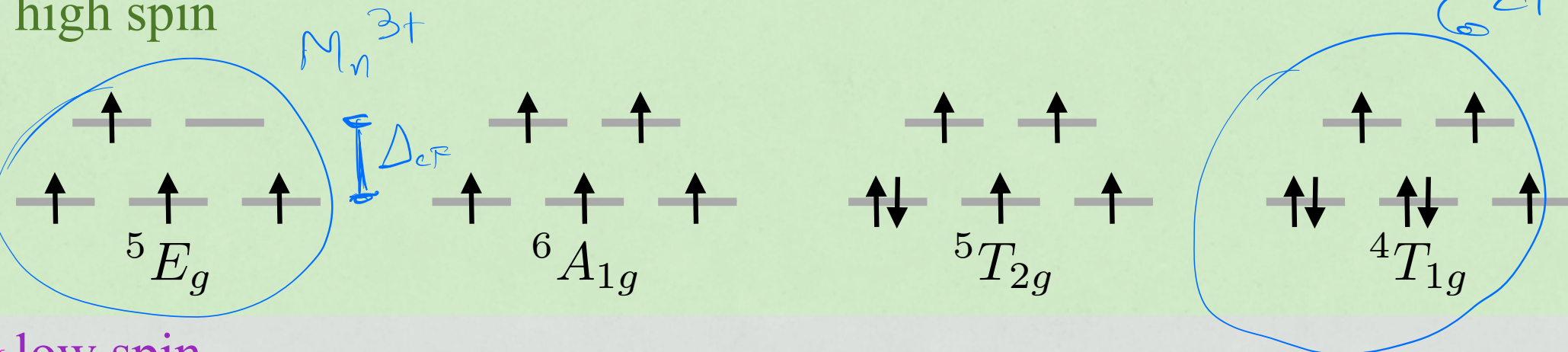
$$E_S - E_T = 2J$$

$$\Rightarrow H_{\text{exch}} = -2J \hat{S}_1 \cdot \hat{S}_2$$

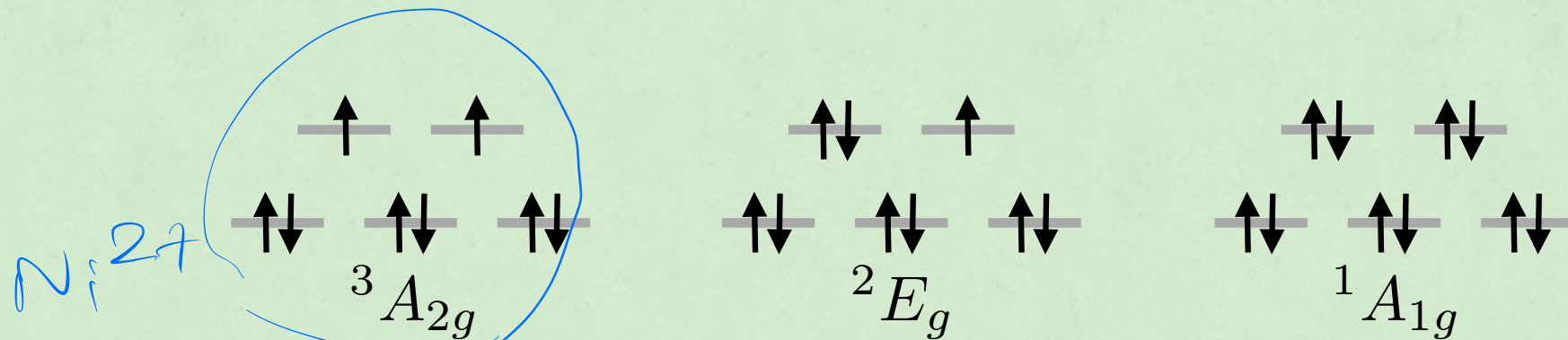
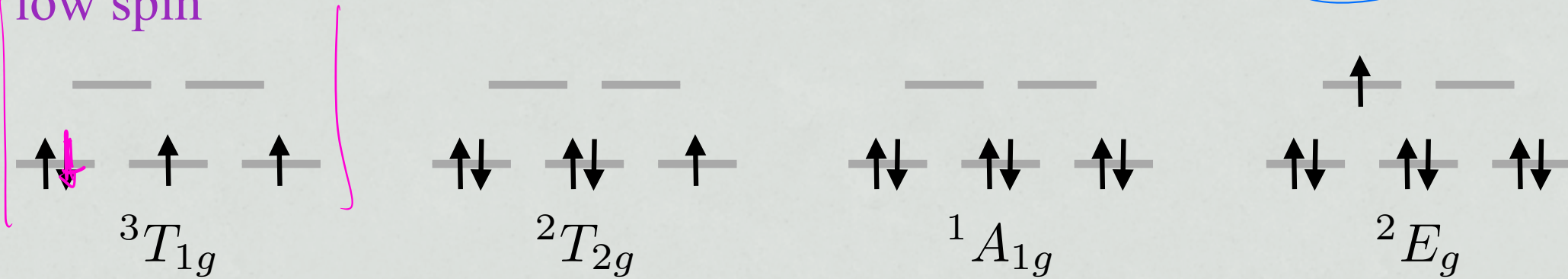
Heisenberg  
effective  
Spin Hamiltonian

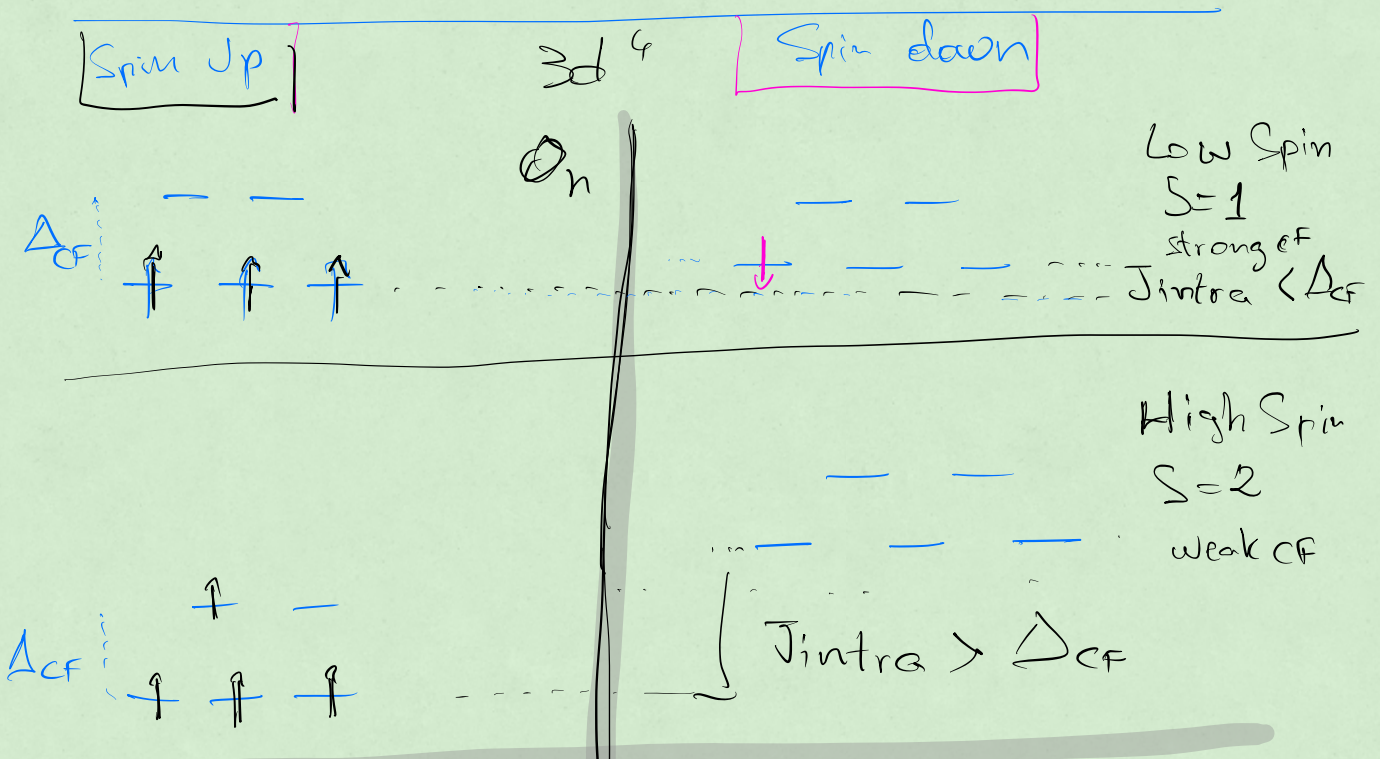
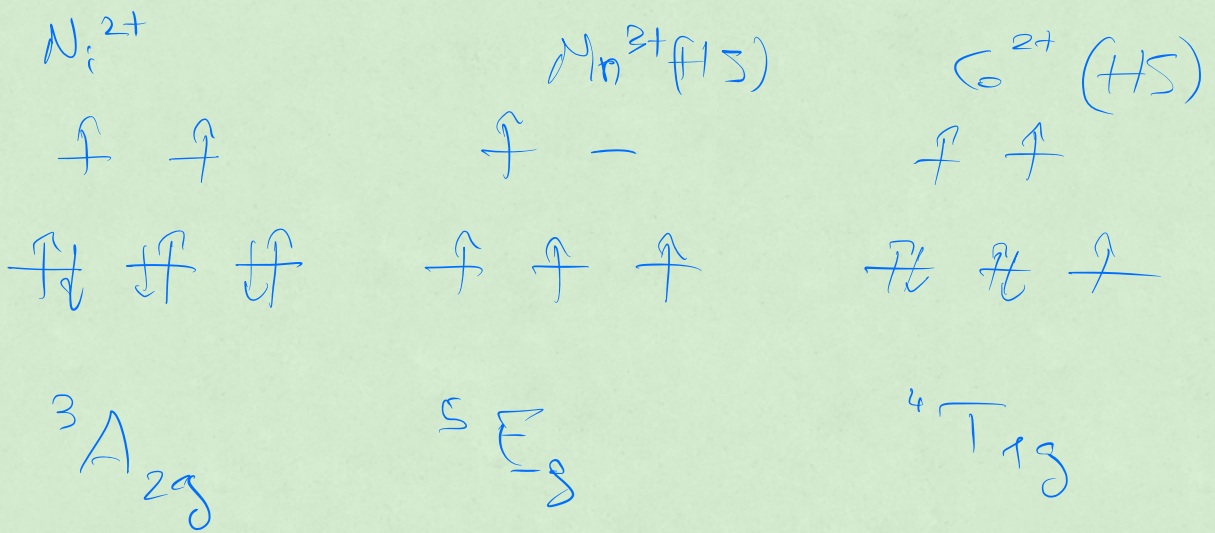


high spin



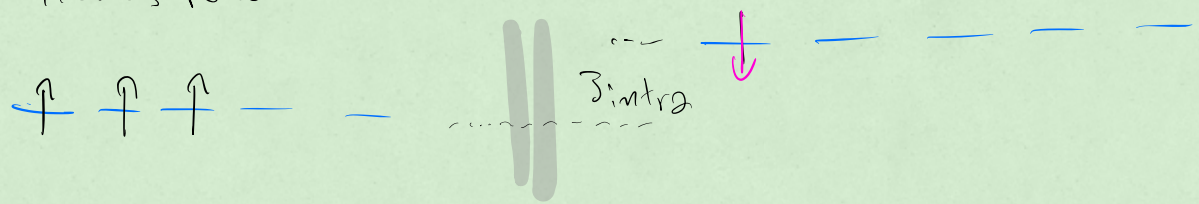
low spin

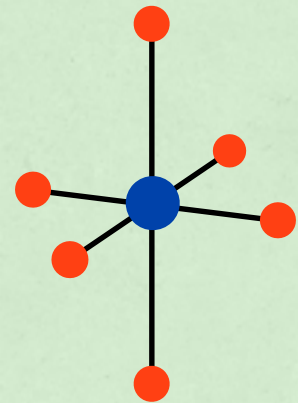




1<sup>st</sup> Hund's rule

Free ion





Elongated

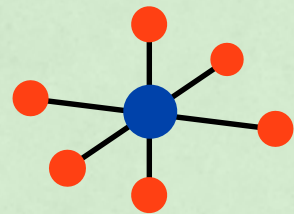
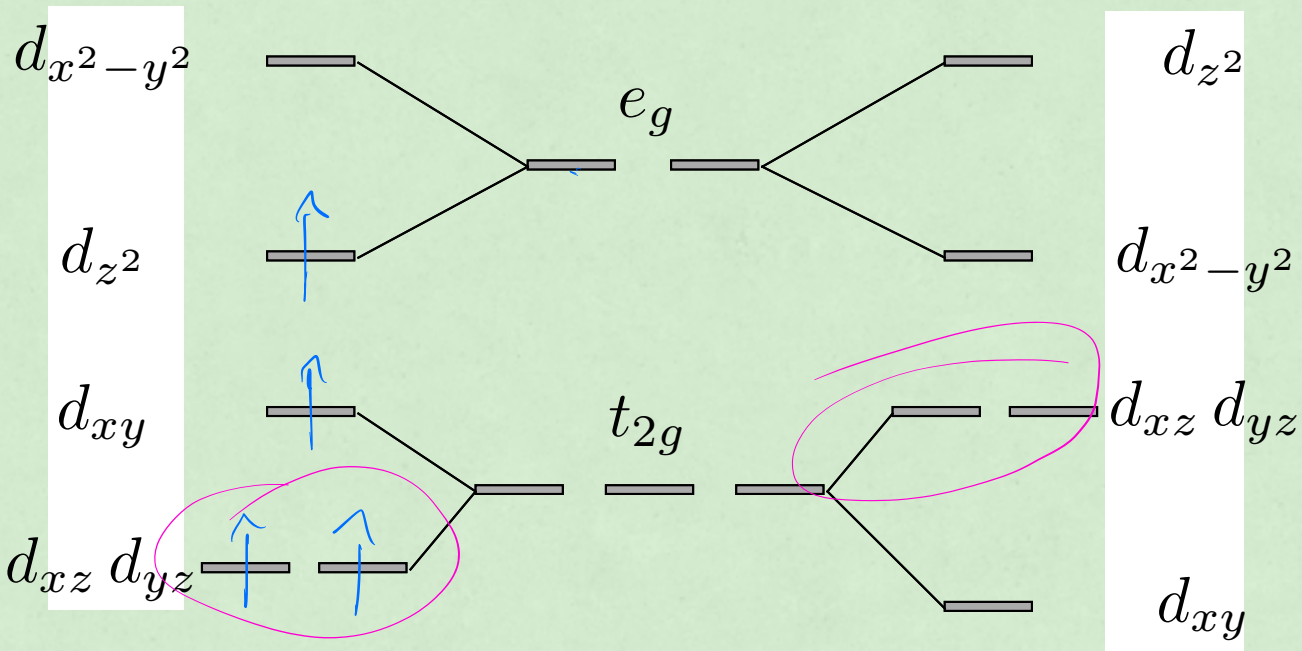
$D_{4h}$

Octahedral

$O_h$

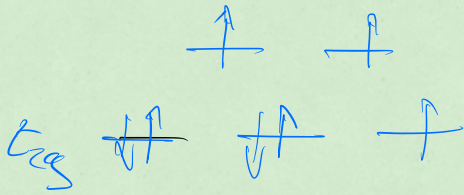
Compressed

$D_{4h}$



degenerate

$G^{2+} (3d)^7$  HS conf



$S = 3/2$   
 $L' = 1$

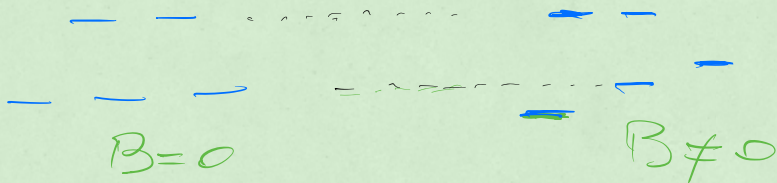
$t_{2g} \begin{cases} d_{xy} \\ d_{yz} \\ d_{xz} \end{cases}$

$H_z = \sum_i H_{z_i} = \sum_i \mu_B B \hat{l}_{z_i}$

$\mu_B B \langle t'_{2g} | \hat{l}_z | t_{2g} \rangle =$

	$d_{xy}$	$d_{xz}$	$d_{yz}$	eigen.	eigst.
$\langle d_{xy}  $	0			0	$ d_{xy}\rangle$
$\langle d_{xz}  $		0	$-i\mu_B B$	$\mu_B B$	$ 2, +1\rangle \propto  d_{xz}\rangle + i d_{yz}\rangle$
$\langle d_{yz}  $		$i\mu_B B$	0	$-\mu_B B$	$ 2, -1\rangle \propto  d_{xz}\rangle - i d_{yz}\rangle$

$\Rightarrow t_{2g}$  behaves like  $l' = 1$  Effectively

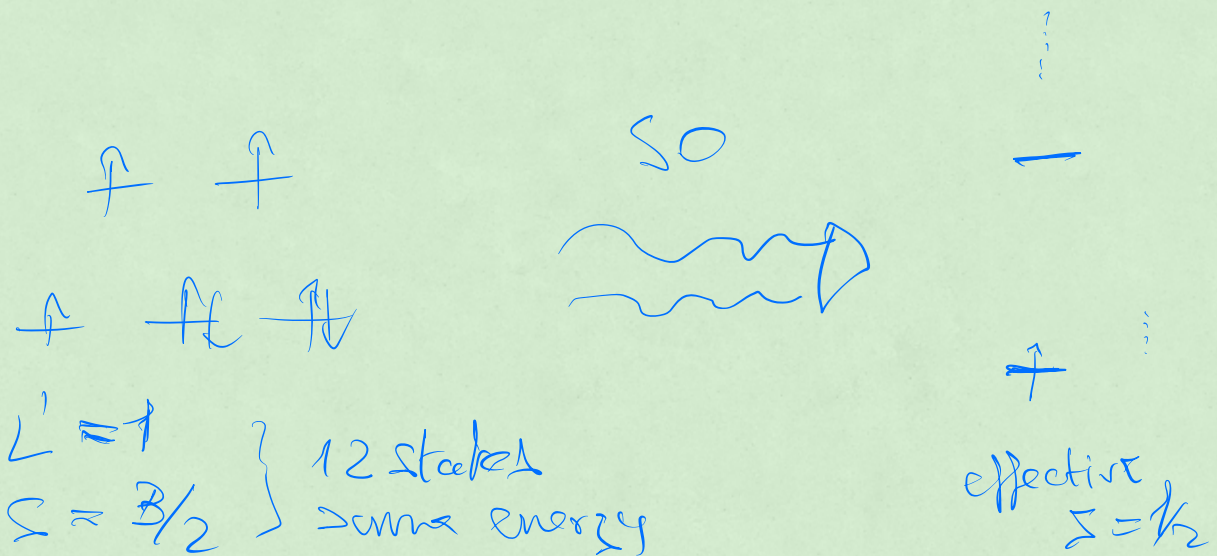


$$\langle e_g | \hat{L}^x | e_g \rangle = 0$$

IS EFFECTIVE

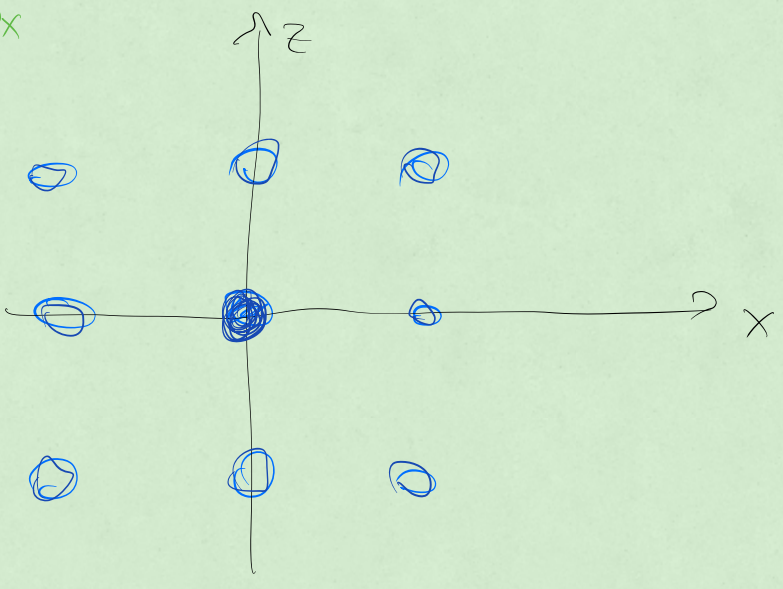
	$ e_g\rangle$	$ e_g\rangle$
1st order	Zee man Splitting	N
	SO coupling	N

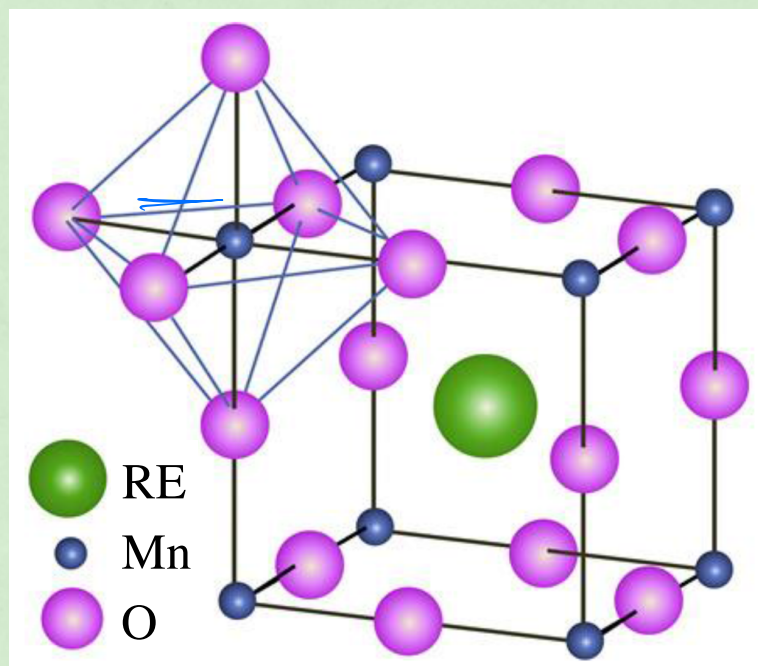
back to  $\text{Co}^{2+} \Rightarrow$  SO is effective at 1st order



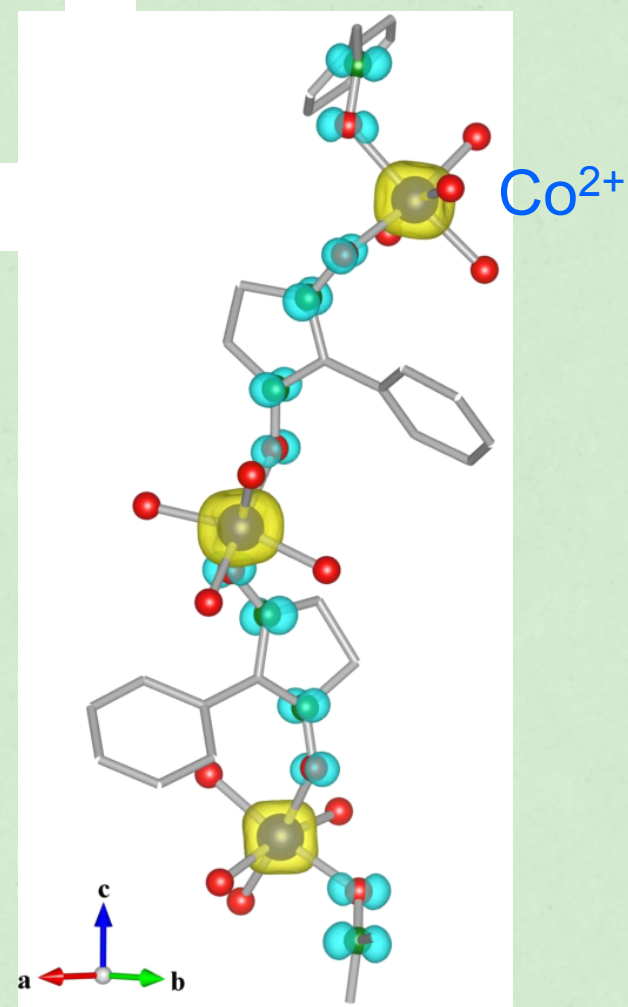


$\omega^{2+}$

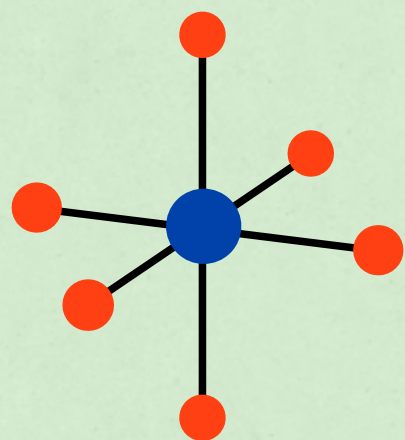




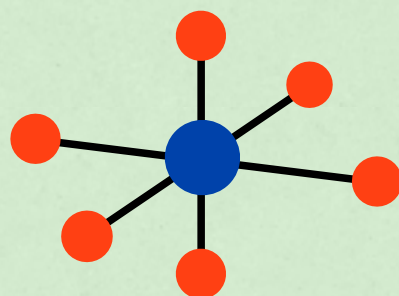
Rare-earth manganites



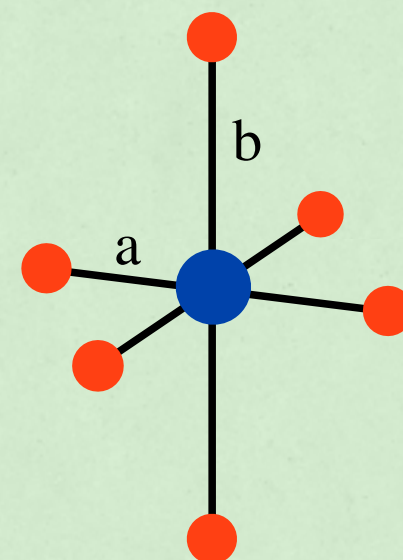
CoPhOMe spin chain



octahedral symmetry  $O_h$



compressed octahedral  
symmetry  $D_{4h}$



elongated octahedral  
symmetry  $D_{4h}$